

**BIDIRECTIONAL REFLECTANCES OF REGOLITHS WITH GRAIN SIZES OF THE ORDER OF THE WAVELENGTH.** B. Hapke<sup>1</sup>, D. DiMucci<sup>1</sup>, R. Nelson<sup>2</sup> and W. Smythe<sup>2</sup>. <sup>1</sup>Dept. Of Geology and Planetary Science, University of Pittsburgh, Pittsburgh, PA 15260. <sup>2</sup>Jet Propulsion Lab., Pasadena, CA 91109.

Scattering and emission of radiation by regoliths with grain sizes  $d$  large compared to the wavelength  $\lambda$ , when the particles scatter quasi-independently, are reasonably well understood. However, this is not true when the particle sizes are comparable with or smaller than  $\lambda$  and interparticle coherent effects can be important, a situation that often occurs in the thermal IR. Nevertheless, many authors treat the two cases similarly. In addition, theoretical calculations [1] have indicated that certain regolith scattering properties, such as the angular width of the opposition effect, should be strongly dependent on  $d/\lambda$ .

In order to understand scattering by particulate media with  $d \sim \lambda$ , we have undertaken a multispectral study of scattering by hematite abrasive powders which span the range from  $d \ll \lambda$  to  $d \gg \lambda$ . Hematite is strongly absorbing at short wavelengths, so that by varying both  $d$  and  $\lambda$  the effects of albedo on the scattering parameters can be separated from the effects of  $d/\lambda$ . The porosities of the powders were all nearly the same, 85-89%.

The powders were viewed by a detector at  $e = 60^\circ$  from the normal and illuminated in the principle plane over a range of angles  $i$  between  $70^\circ$  on both sides of the normal. The phase angle  $g$  varied between  $1^\circ$  and  $130^\circ$ . The data were then fitted by an equation of the form (2)

$$r(i, e, g) = (w/4) * \mu_o / (\mu_o + \mu) * OE(g) * [p(g) + H(\mu_o)H(\mu) - 1],$$

$$OE(g) = 1 + A / [1 + (1/h) \tan(g/2)],$$

$$p(g) = (1+c)/2 * (1-b^2) / (1 - 2b \cos g + b^2)^{3/2} + (1-c)/2 * (1b^2) / (1 + 2b \cos g + b^2)^{3/2},$$

$$H(x) = [1 + 2x] / [1 + 2x * \text{sqr}(x)],$$

where  $r$  = bidirectional reflectance,  $w$  = single scattering albedo,  $\mu_o = \cos i$ ,  $\mu = \cos e$ . The quantities  $w$ ,  $A$ ,  $h$ ,  $b$  and  $c$  are the parameters that were fitted. The particle phase function  $p(g)$  is a double Henyey-Greenstein function;  $c$  determines whether the particle

is forward ( $c > 0$ ) or back ( $c < 0$ ) scattering and  $b$  determines the width and amplitude of the lobes.  $OE(g)$  is an opposition effect function with amplitude  $A$  and angular width  $h$  that includes the averaged effects of both a shadow hiding and a coherent backscatter opposition effect.

When an isolated particle becomes smaller than the wavelength its scattering efficiency decreases and its phase function is proportional to  $1 + \cos^2 g$ . Hence, if the particles in a powder behaved similarly, it would be expected that when  $d/\lambda < 1$ ,  $w$  would decrease markedly and  $c \rightarrow 0$ . This is not what is observed, as shown in Figures 1 and 2. All sizes of particles were strongly back scattering ( $c > 0$ ) and the single scattering albedo increased, rather than decreased for small particles.

If the opposition effect is caused primarily by coherent backscatter,  $h$  would be expected to decrease as  $d/\lambda$  increases, while  $A$  should increase with  $w$ . Figure 3 shows that  $h$  is not correlated with  $d/\lambda$ , which is what would be expected if a large part of the opposition effect is caused by shadow hiding.

These results show that particles smaller than the wavelength in a regolith do not scatter independently. Apparently, coherent interactions between them cause clusters of small particles to scatter like much larger particles. Although coherent backscatter does occur, as indicated in Figure 4 by the increase of  $A$  with  $w$  (and also by an upturn in the circular polarization ratio at small  $g$ , which was also measured, but is not shown), these clusters apparently dominate the scattering and cast shadows that cause a shadow hiding opposition effect.

**REFERENCES:** [1] M. Mishchenko (1992), *Astrophys. Space Sci* **194**, 327. [2] B. Hapke (1993), *Theory of Reflectance and Emittance Spectroscopy*, Cambridge U. Press, N. Y.

